



Energy Methods For Initial –Boundary String Problem

Aqeel F. Jaddoa

Dept. of Mathematics/College of Education for Pure Science (Ibn Al-Haitham)/University of Baghdad

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Abstract

We study one example of hyperbolic problems it's Initial-boundary string problem with two ends. In fact we look for the solution in weak sense in some sobolev spaces. Also we use energy technic with Galerkin's method to study some properties for our problem as existence and uniqueness.

Keyword: Initial- Boundary String Problem, Weak Solution, Sobolev Spaces, Galerkin's Method, Energy Estimate.



1. Introduction

Throughout this paper we consider energy estimate for string equation with initial –boundary conditions. We gave a little introduction for weak derivatives, weak solution in some sobolev space, for example $H_0^1(u)$, $H^{-1}(u)$ and others. In fact energy estimate has many results for example a speed of propagation [1], in our paper it leads us to straight forward existence, uniqueness [2], results and other properties like regularity. It yields also to well-posed for hyperbolic equations.

In 2007, Tarama [3], introduced wave equations with coefficients satisfying Besov type conditions and obtained the energy estimate, also he gave an example of a wave equation with continuous and differentiable coefficients for which the L^2 estimate holds.

Jaipong Kasemsuwan [4], concerned with the energy decay of the global solution for initial-boundary value problem to a nonlinear damped equation of suspended string with uniform density to which a nonlinear outer force works.

The string equation with initial-boundary conditions is

$$\frac{\partial^2 u}{\partial t^2} = a^2 \frac{\partial^2 u}{\partial x^2} + f(x, t) \text{ in } Q \quad \dots (1.1)[5, p.30]$$

Where $Q = \{(x, t); x \in (0, l), t \in (0, T)\}$,

Initial conditions:

$$\begin{cases} u(x, 0) = \varphi(x) \\ u_t(x, 0) = \psi(x) \end{cases} \dots (1.2)$$

Boundary conditions:

$$\begin{cases} u(0, t) = 0 \\ u(l, t) = 0 \end{cases} \dots (1.3)$$

This problem describes the vibration of the string in a finite interval $[0, l]$ and the constant a denotes the speed of the string.

Rewrite the initial-boundary problem with the second order partial differential operator L as:

$$\left. \begin{array}{l} \frac{\partial^2 u}{\partial t^2} = a^2 Lu + f(x, t) \text{ in } U \times (0, T) \\ u(x, t) = 0 \text{ on } \partial U \times [0, T] \\ u(x, t) = \varphi(x), \quad \frac{\partial}{\partial t} u(x, t) = \psi(x) \text{ on } U \times \{t = 0\} \end{array} \right\} \dots (1.4)$$

Let us denote for $U = (0, l)$, $U_T = U \times (0, T]$, $f: U \times (0, T] \rightarrow \mathcal{R}$ and $\varphi, \psi: U \rightarrow \mathcal{R}$ and these functions are given.

The function $u: \overline{U_T} \rightarrow \mathcal{R}$ is the unknown function $u=u(x, t)$. The symbol L denotes for each time t a second-order partial differential operator, having the divergence form

$$Lu = - \sum_{i,j=1}^n (a^{ij}(x, t) u_{xi})_{xj} + \sum_{i=1}^n b^i(x, t) u_{xi} + c(x, t) u \quad \dots (1.5) [6, p. 309]$$

For given coefficients $a^{ij}, b^i, c \in C^1(\overline{U} \times [0, T])$ ($i, j = 1, 2, \dots, n$), $f \in L^2(U \times [0, T])$, $\varphi \in H_0^1(U)$,

$\psi \in L^2(U)$, And always assume that $a^{ij} = a^{ji}$ ($i, j = 1, 2, \dots, n$).

Definition 1[6, p.378]

The partial differential operator $\frac{\partial^2}{\partial t^2} + L$ is said to be (uniformly) hyperbolic if there exists a constant $\theta > 0$ such that $\sum_{i,j=1}^n a^{ij}(x, t) \varepsilon_i \varepsilon_j \geq \theta |\varepsilon|^2$ for all $(x, t) \in U \times [0, T]$, $\varepsilon \in R^n$.

Definition 2 [6, p. 285]

Let $u \in L^1(0, T, X)$. we say $v \in L^1(0, T, X)$ is **weak derivative** of u , written: $u' = v$ provided



$$\int_0^T \varphi'(t) u(t) dt = - \int_0^T \varphi(t) v(t) dt$$

For all scalar test functions $\varphi \in C_c^\infty(0, T)$.

Let us introduce the time-dependent bilinear form:

$$B[u, v; t] = \int_U \left(\sum_{i,j=1}^n a^{ij} u_{xi} v_{xj} \right) dx + \int_U \sum_{i=1}^n (b^i(x, t) u_x v + c(x, t) u v) dx$$

for all $u, v \in H_0^1(U)$ and $0 \leq t \leq T$.

First we suppose u is smooth solution and defined the mapping:

$$\begin{aligned} u: [0, T] &\rightarrow H_0^1(U), [u(t)](x) = u(x, t), (x \in U, 0 \leq t \leq T), (') = \frac{d}{dt} \\ f: [0, T] &\rightarrow L^2(U), [f(t)](x) = f(x, t) \end{aligned}$$

Now we fix any function $v \in H_0^1(U)$, multiply by v the PDE $(u_{tt} + Lu = f)v$

$$(u'', v) + B[u, v, t] = (f, v)$$

The pairing $(,)$ inner product in $L^2(U)$.

We see the form: $u_{tt} = \varphi^0 + \sum_{j=1}^n \varphi_{xj}^j$

$$\text{For } \varphi^0 = f - \sum_{i=1}^n b^i u_{xi} - c u, \quad \varphi^j = f - \sum_{i=1}^n a^{ij} u_{xi} \quad (j=1, \dots, n)$$

This suggests that we should look for weak solution u with $u'' \in H^{-1}(U)$ for almost every where $0 \leq t \leq T$, $\langle u'', v \rangle$ means $u'' \in H^{-1}(U), v \in H_0^1(U)$

Definition 3 [6, p.379]

We say a function $u \in L^2(0, T, H_0^1(U))$, with $u' \in L^2(0, T, L^2(U))$, $u'' \in L^2(0, T, H^{-1}(U))$

is **weak solution** of the hyperbolic initial(boundary) problem and

$$\text{i. } \langle u'', v \rangle + B[u, v, t] = (f, v) \text{ for each } v \in H_0^1(U)$$

$$\text{ii. } u(0) = \varphi, \quad u'(0) = \psi.$$

As known sobolev spaces are designed to contain less smooth functions that means the functions have some but not great, smoothness properties.

So the idea is to make approximation for the function $u(x, t)$ such that $u(x, t)=0$ in the boundary of U by the

$$u_m(t) = \sum_{k=1}^m d_m^k(t) w_k(x) \quad [6, p. 353]$$

In the right side hand we have separated variables, we note that U is a domain of the variable x only, and this is the reason to consider the variable t in respect to the sobolev space $H_0^1(U)$ as a parameter. where $w_k = w_k(x)$ ($k = 1, 2, \dots$) are smooth functions such that $\{w_k\}_{k=1}^\infty$ is an orthogonal basis of $H_0^1(U)$, $\{w_k\}_{k=1}^\infty$ is an orthonormal basis of $L^2(U)$.

We are going to consider u not as a function of x and t together, but rather as a mapping u of t into the space $H_0^1(U)$ of functions of x .

We will look for a function $u_m: [0, T] \rightarrow H_0^1(U)$ of the form $u_m(t) = \sum_{k=1}^m d_m^k(t) w_k$ such that

$$(u_m'', w_k) + B[u_m, w_k; t] = (f, w_k), \quad (t \in [0, T], k = 1, 2, \dots, m) \quad \dots (1.6)$$

$$\begin{cases} d_m^k(0) = (\varphi, w_k) \\ d_m^k'(0) = (\psi, w_k) \end{cases} \quad \dots (1.7)$$

where

$$B[u_m, w_k; t] = \int_U \left(\sum_{i,j=1}^n a^{ij} u_{m,xi} w_{k,xj} \right) dx + \int_U \sum_{i=1}^n (b^i u_{m,xi} w_k + c u_m w_k) dx$$

In other words, we construct solutions of certain finite-dimensional approximations to our problem and then passing to limits, it's called Galerkin's method [6, p.380].

Indeed we need some estimates to send m to infinity and to show subsequence of our solutions u_m of the approximate problem .

2. Main Result

Theorem (Energy estimates)(2.1): *There exists a constant C , depending only on U, T and the coefficients on L such that:*

$$\max_{0 \leq t \leq T} (\|u_m(t)\|_{H_0^1(U)} + \|u'_m(t)\|_{L^2(U)}) + \|u''_m(t)\|_{L^2(0,T;H^{-1}(U))} \\ \leq C(\|f\|_{L^2(0,T;L^2(U))} + \|\varphi\|_{H_0^1(U)} + \|\psi\|_{L^2(U)}) \quad \text{for } m=1,2,\dots$$

proof We take equation (1.6) and write it in details as:

$$\int_U (u''_m w_k) dx + \int_U \sum_{i,j=1}^n (a^{ij} u_{m,xi} w_{k,xj}) dx + \int_U \sum_{i=1}^n (b^i u_{m,xi} w_k + c u_m w_k) dx \\ = \int_U f w_k dx$$

Multiply by $d_m^{k'}(t)$ so:

$$\int_U [u''_m w_k d_m^{k'}(t)] dx + \int_U \left(\sum_{i,j=1}^n a^{ij} u_{m,xi} w_{k,xj} d_m^{k'}(t) + \sum_{i=1}^n b^i u_{m,xi} w_k d_m^{k'}(t) \right. \\ \left. + c u_m w_k d_m^{k'}(t) \right) dx = \int_U f w_k d_m^{k'}(t) dx.$$

$$\int_U (u''_m u'_m) dx + \int_U \left(\sum_{i,j=1}^n a^{ij} u_{m,xi} u'_{m,xj} + \sum_{i=1}^n b^i u_{m,xi} u'_m + c u_m u'_m \right) dx = \int_U f u'_m dx \\ (u''_m, u'_m) + B[u_m, u'_m; t] = (f, u'_m).$$

And we have

$$(u''_m, u'_m) = \frac{d}{dt} \left(\frac{1}{2} \|u'_m\|_{L^2(U)}^2 \right).$$

We write :

$$\int_U \sum_{i,j=1}^n (a^{ij} u_{m,xi} u'_{m,xj}) dx + \int_U \sum_{i=1}^n (b^i u_{m,xi} u'_m + c u_m u'_m) dx = B_1 + B_2$$

where

$$B_1 = \int_U \left(\sum_{i,j=1}^n a^{ij} u_{m,xi} u'_{m,xj} \right) dx \quad \text{and} \quad B_2 = \int_U \left(\sum_{i=1}^n b^i u_{m,xi} u'_m + c u_m u'_m \right) dx$$

Now we define the symmetric bilinear form: $A[u, v; t] := \int_U (\sum_{i,j=1}^n a^{ij} u_{xi} v_{xj}) dx$,

$$u, v \in H_0^1(U)$$

Then

$$\frac{d}{dt} \left(\frac{1}{2} A[u_m, u_m; t] \right) - \frac{1}{2} \int_U \left(\sum_{i,j=1}^n a^{ij} u_{m,xi} u_{m,xj} \right) dx \\ = \frac{1}{2} \left(\int_U \frac{d}{dt} \left[\sum_{i,j=1}^n a^{ij} u_{m,xi} u_{m,xj} \right] dx - \frac{1}{2} \int_U \left(\sum_{i,j=1}^n a_t^{ij} u_{m,xi} u_{m,xj} \right) dx \right)$$

$$\begin{aligned}
 &= \frac{1}{2} \int_U \sum_{i,j=1}^n a_t^{ij} u_{m_{xi}} u_{m_{xj}} dx + \frac{1}{2} \int_U \sum_{i,j=1}^n a^{ij} u'_{m_{xi}} u_{m_{xj}} dx + \frac{1}{2} \int_U \sum_{i,j=1}^n a^{ij} u'_{m_{xi}} u_{m_{xj}} dx \\
 &- \frac{1}{2} \int_U \sum_{i,j=1}^n a_t^{ij} u_{m_{xi}} u_{m_{xj}} dx = \int_U \sum_{i,j=1}^n a^{ij} u_{m_{xi}} u'_{m_{xj}} dx = B_1
 \end{aligned}$$

(since $a^{ij} = a^{ji}$)

Then

$$B_1 = \frac{d}{dt} \left(\frac{1}{2} A[u_m, u_m; t] \right) - \frac{1}{2} \int_U \sum_{i,j=1}^n a_t^{ij} u_{m_{xi}} u_{m_{xj}} dx$$

So

$$\frac{d}{dt} \left(\frac{1}{2} A[u_m, u_m; t] \right) = B_1 + \frac{1}{2} \int_U \left(\sum_{i,j=1}^n a_t^{ij} u_{m_{xi}} u_{m_{xj}} \right) dx$$

$a_t^{ij} \in C^1(\bar{U} \times [0, T])$, so there exists $\text{Max}(a_t^{ij})$ such that

$$\begin{aligned}
 B_1 + \frac{1}{2} \int_U \left(\sum_{i,j=1}^n a_t^{ij} u_{m_{xi}} u_{m_{xj}} \right) dx &\leq B_1 + C \int_U \left(\sum_{i,j=1}^n u_{m_{xi}} u_{m_{xj}} \right) dx \\
 &\leq B_1 + C \int_U \sum_{j=1}^n u_{m_{xj}}^2 dx \\
 &\leq B_1 + C \|u_m\|_{H_0^1}^2
 \end{aligned}$$

Or we can write it as:

$$\frac{d}{dt} (A[u_m, u_m; t]) = 2B_1 + C \|u_m\|_{H_0^1}^2 \quad \dots (2.1)$$

$$\text{We had } (u''_m, u'_m) = \frac{d}{dt} \left(\frac{1}{2} \|u'_m\|_{L^2(U)}^2 \right)$$

So

$$\begin{aligned}
 \frac{d}{dt} (\|u'_m\|_{L^2(U)}^2) &= 2(u''_m, u'_m) = 2(f, u'_m) - 2(B_1 + B_2) \\
 &\leq C(\|f\|_{L^2}^2 + C \|u'_m\|_{L^2}^2) - 2(B_1 + B_2)
 \end{aligned}$$

Now we obtain

$$\frac{d}{dt} (\|u'_m\|_{L^2(U)}^2) \leq \|f\|_{L^2}^2 + C \|u'_m\|_{L^2}^2 - 2(B_1 + B_2) \quad \dots (2.2)$$

But

$$B_2 = \int_U \left(\sum_{i=1}^n b^i u_{m_{xi}} u'_m \right) dx \rightarrow |B_2| = \left| \int_U \left(\sum_{i=1}^n b^i u_{m_{xi}} u'_m \right) dx \right|$$

$b^i \in C^1(\bar{U} \times [0, T])$ then there exists a Max (b^i) such that

$$\begin{aligned}
 &\left| \int_U \sum_{i=1}^n (b^i u_{m_{xi}} u'_m + c u_m u'_m) dx \right| \leq C \int_U \sum_{i=1}^n (|u_{m_{xi}} u'_m| + |u_m u'_m|) dx \\
 &\leq C \int_U \sum_{i=1}^n (u_{m_{xi}}^2 + u_m^2 + 2u'^2_m) dx
 \end{aligned}$$



$$\leq C (\|u_m\|_{H_0^1}^2 + \|u'_m\|_{L^2(U)}^2)$$

Or

$$|B_2| \leq C (\|u_m\|_{H_0^1(U)}^2 + \|u'_m\|_{L^2(U)}^2)$$

From (2.1) and (2.2) we obtain :

$$\begin{aligned} \frac{d}{dt} (A[u_m, u_m; t] + \|u'_m\|_{L^2(U)}^2) \\ \leq 2B_1 + C\|u_m\|_{H_0^1(U)}^2 + C\|f\|_{L^2(U)}^2 + C\|u'_m\|_{L^2(U)}^2 - 2(B_1 + B_2) \\ \leq C\|u_m\|_{H_0^1(U)}^2 + C\|f\|_{L^2(U)}^2 + C\|u'_m\|_{L^2(U)}^2 + C|B_2| \\ \leq C\|u_m\|_{H_0^1(U)}^2 + C\|u'_m\|_{L^2(U)}^2 + C\|f\|_{L^2(U)}^2 + C\|u_m\|_{L^2}^2 \\ \leq C\|u_m\|_{H_0^1(U)}^2 + C\|u'_m\|_{L^2(U)}^2 + C\|f\|_{L^2(U)}^2 \end{aligned}$$

So we have

$$\frac{d}{dt} (A[u_m, u_m; t] + \|u'_m\|_{L^2(U)}^2) \leq C\|u_m\|_{H_0^1(U)}^2 + C\|u'_m\|_{L^2(U)}^2 + C\|f\|_{L^2(U)}^2$$

But L is hyperbolic operator and by the definition of hyperbolic operator we have

$$\theta \int_U \sum_{i=1}^n u_{mxi} dx \leq \int_U \sum_{i,j=1}^n a^{ij} u_{mxi} u_{mxj} dx = A[u_m, u_m; t].$$

$$\theta \int_U \sum_{i=1}^n u_{mxi} dx \leq A[u_m, u_m; t].$$

And

$$C\|u_m\|_{H_0^1(U)}^2 \leq C \int_U u_m^2 dx + C \int_U \sum_i^n u_{mxi}^2 dx \leq A[u_m, u_m; t]$$

$$\frac{d}{dt} (A[u_m, u_m; t] + \|u'_m\|_{L^2(U)}^2) \leq C\|u'_m\|_{L^2(U)}^2 + C\|f\|_{L^2(U)}^2 + C A[u_m, u_m; t]$$

We denote

$$\eta(t) = \|u'_m(t)\|_{L^2(U)}^2 + A[u_m(t), u_m(t); t]$$

$$\xi(t) = \|f\|_{L^2(U)}^2$$

$$So we have \eta'(t) \leq C_1 \eta(t) + C_2 \xi(t)$$

Now we use Gronwall Inequality

$$\eta(t) \leq e^{c_1 t} (\eta(0) + c_2 \int_0^t \xi(s) ds) \quad ... (2.3)$$

But

$$\begin{aligned} \eta(0) &= \|u'_m(0)\|_{L^2(U)}^2 + A[u_m(0), u_m(0); 0] \\ &\leq \|u'_m(0)\|_{L^2(U)}^2 + \int_U \sum_{i,j=1}^n a^{ij} u_m(0)_{xi} u_m(0)_{xj} dx \\ &\leq \|u'_m(0)\|_{L^2(U)}^2 + \int_U \left(\sum_{i,j=1}^n a^{ij} u_m(0)_{xi} u_m(0)_{xj} \right) dx = \\ &= \|u'_m(0)\|_{L^2(U)}^2 + \|u_m(0)\|_{H_0^1(U)}^2 \\ &\leq C(\|\psi\|_{L^2(U)}^2 + \|\varphi\|_{H_0^1(U)}^2) \end{aligned}$$

so



$$\eta(0) \leq C(\|\psi\|_{L^2(U)}^2 + \|\varphi\|_{H_0^1(U)}^2)$$

from (2.3) we will get

$$\|u'_m(t)\|_{L^2(U)}^2 + A[u_m(t), u_m(t); t] \leq C \left(\|\psi\|_{L^2(U)}^2 + \|\varphi\|_{H_0^1(U)}^2 + \int_0^t \|f\|_{L^2(0,T;L^2(U))}^2 dt \right)$$

so

$$\|u'_m(t)\|_{L^2(U)}^2 + A[u_m(t), u_m(t); t] \leq C \left(\|\psi\|_{L^2(U)}^2 + \|\varphi\|_{H_0^1(U)}^2 + \|f\|_{L^2(0,T;L^2(U))}^2 \right).$$

But again L is hyperbolic (2.3)

$$so \|u'_m(t)\|_{L^2(U)}^2 + \|u_m(t)\|_{H_0^1(U)}^2 \leq C \left(\|\psi\|_{L^2(U)}^2 + \|\varphi\|_{H_0^1(U)}^2 + \|f\|_{L^2(0,T;L^2(U))}^2 \right)$$

but t is an arbitrary

so

$$\begin{aligned} \max_{0 \leq t \leq T} (\|u'_m(t)\|_{L^2(U)}^2 + \|u_m(t)\|_{H_0^1(U)}^2) &\leq \\ C \left(\|\psi\|_{L^2(U)}^2 + \|\varphi\|_{H_0^1(U)}^2 + \|f\|_{L^2(0,T;L^2(U))}^2 \right) &\dots (2.4) \end{aligned}$$

Now Let $v \in H_0^1(U)$ and $\|v\|_{H_0^1(U)} \leq 1$.

We write $v = v_1 + v_2$, where $v_1 \in \text{span}\{w_k\}_{k=1}^m$ and $(v_2, w_k) = 0$.
We have

$$\begin{aligned} \langle u''_m, v \rangle &= (u''_m, v) = \left(\sum_{k=1}^m w_k d_m^k''(t), v_1 + v_2 \right) \\ &= \left(\sum_{k=1}^m w_k d_m^k''(t), v_1 \right) + \left(\sum_{k=1}^m w_k d_m^k''(t), v_2 \right) \\ &\quad (u''_m, v_1) = (f, v_1) - B[u_m, v, ; t] \end{aligned}$$

but

$$\begin{aligned} (f, v_1) &\leq \|f\|_{L^2(0,T;L^2(U))} \|v_1\|_{L^2(U)} \\ &\leq \|f\|_{L^2(0,T;L^2(U))} \end{aligned}$$

So

$$B[u_m, v; t] \leq \|u_m(t)\|_{H_0^1(U)} \|v_1\|_{H_0^1(U)}$$

$$\leq \|u_m(t)\|_{H_0^1(U)}.$$

$$|B| \leq C \|u_m(t)\|_{H_0^1(U)}$$

and we have

$$|(f, v_1)| \leq \|f\|_{L^2(0,T;L^2(U))}$$

So

$$\begin{aligned} | \langle u''_m, v \rangle | &\leq |(f, v_1)| + |B| \\ &\leq C(\|f\|_{L^2(0,T;L^2(U))} + \|u_m(t)\|_{H_0^1(U)}) \end{aligned}$$

By the definition of

$$\|u''_m u(t)\|_{H^{-1}(U)} = \max_{\|v\| \leq 1} | \langle u''_m, v \rangle |.$$

$$\|u''_m\|_{H^{-1}(U)} \leq C(\|f\|_{L^2(0,T;L^2(U))} + \|u_m(t)\|_{H_0^1(U)}).$$

So

$$\int_0^T \|u''_m\|_{H^{-1}(U)} dt \leq C \int_0^T (\|f\|_{L^2(0,T;L^2(U))}^2 + \|u_m(t)\|_{H_0^1(U)}^2) dt$$

$$\leq C \|f\|_{L^2(0,T;L^2(U))}^2 + \int_0^T \|u_m(t)\|_{H_0^1(U)}^2 dt$$



$$\begin{aligned}
 &\leq C \|f\|_{L^2(0,T;L^2(U))} + \max \|u_m(t)\|_{H_0^1(U)}^2 \int_0^T dt \\
 &\quad \text{From (2.4)} \\
 &\leq C \left(\|\varphi\|_{H_0^1(U)}^2 + \|\psi\|_{L^2(0,T;L^2(U))}^2 + \|f\|_{L^2(0,T;L^2(U))}^2 \right). \\
 &\quad \text{So} \\
 \|u_m''\|_{L^2(0,T,H^{-1})} &\leq C \left(\|\varphi\|_{H_0^1(U)}^2 + \|\psi\|_{L^2(0,T;L^2(U))}^2 + \|f\|_{L^2(0,T;L^2(U))}^2 \right) \quad \dots \\
 &\quad (2.5) \\
 (2.4) + (2.5) &= (2.3)
 \end{aligned}$$

Conclusion

In our main result, we gave a good estimate to measure the regularity by using some Sobolev spaces(integral spaces) which give the property of continuity to the function u which is the solution of our problem and the continuity of first derivative, second derivative more regular of course in integral sense. Also, $u=0$ on the boundary of the domain U is generalized sense, this means when we deal with an integral space we say that the trace of u is equal to zero on the boundary of the domain U .

References

1. Hirosawa, F., (2007), On the Asymptotic Behavior of the Energy for the Wave Equations with Time Depending Coefficients, *Math. Ann.*, 339, 819-839.
2. Faciu ,C. and Simion, N. (2000), Energy Estimates and Uniqueness of the Weak Solutions of Initial – Boundary Value Problems for Semilinear Hyperbolic Systems, *Z. Angew. Math.*, 51, 792-805.
3. Shigeo, T. (2007), Energy Estimate for Wave Equations with Coefficients in Some Besov Type Class, *Electronic Journal of Equations*, 2007(85)1-12.
4. Jaipong, K. (2011), Exponential Decay for Nonlinear Damped Equation of Suspended StringProc. Of CSIT, 1, 309-313.
5. Michael, R. and Robert, C. (2000), An Introduction to Partial Differential Equations, Springer.
6. Lawrence, C. Evans, (1997), Partial Differential Equations", American Mathematical Society.



طرائق الطاقة لمسألة الوتر ذات القيم الابتدائية والحدودية

عقيل فالح جدع

قسم الرياضيات / كلية التربية للعلوم الصرفة (أبن الهيثم) / جامعة بغداد

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الخلاصة

درسنا أحد أمثلة مسائل القطوع الزائدة وهي مسألة الوتر ذو النهايتين ذات القيم الابتدائية والحدودية. في الحقيقة نحن نبحث عن حل بمعنى الحل الضعيف لبعض فضاءات سوبولف واستخدمنا طرائق الطاقة مع طريقة Galerkin) أيضاً دراسة بعض الخصائص لمسألتنا مثل الوجود والوحدانية.

الكلمات المفتاحية: مسألة الوتر ذات القيم الابتدائية والحدودية، الحل الضعيف، فضاءات سوبولف، طريقة Galerkin، تخمين الطاقة